

Steepening of Alfvén waves due to parallel compressibility

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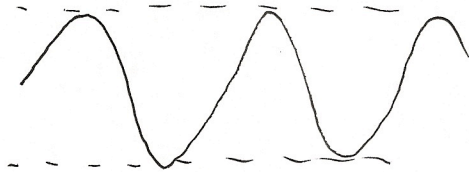
(Ref: "Relaxation Dynamics in Laboratory and Astrophysical Plasmas"
Ch 4.4, P.H. Diamond, et. al. 2010)

Alfvén waves $\leftrightarrow \nabla \cdot \underline{v} = 0$ ~~→~~ Steepening

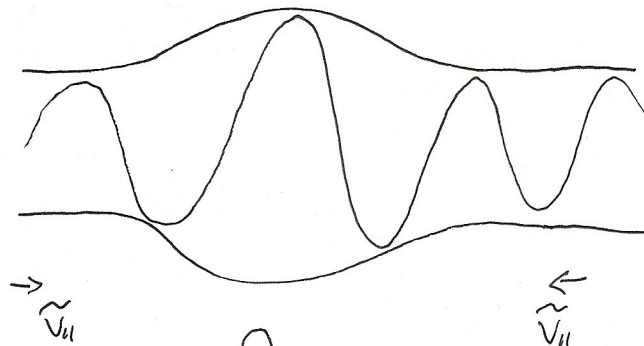
However, a little compressibility will lead to steepening of wave ~~packet~~ ^{packet}.

Approach: Modulational instability

Alfvén wave:

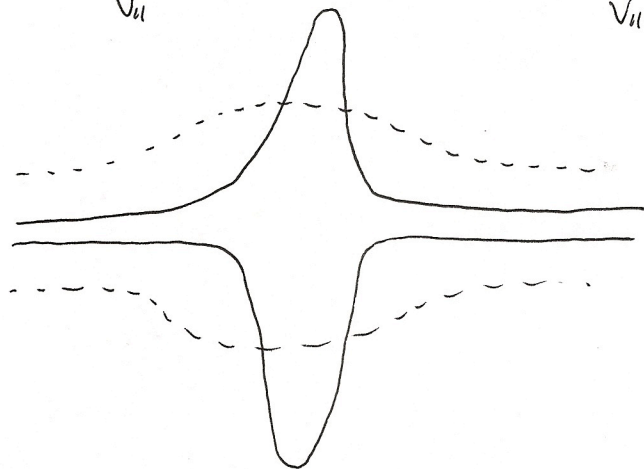


Small compression:



\Rightarrow Parallel flow $\tilde{v}_{||}$
 $\tilde{v}_{||} \sim \tilde{B}^2$
 \downarrow
 $\tilde{B} \uparrow, \frac{\partial}{\partial t} \tilde{B} = -\underline{v} \cdot \nabla \tilde{B}$
feedback loop

Steepening:



Result.

Until: Saturated by

- dissipation η
- dispersion $d_i \sim \frac{c}{\omega_{pi}} \leftarrow$ ion inertial scale

Physical derivation :

②

Need equation for envelope. \leftrightarrow Scale separation $\begin{matrix} \rightarrow \text{fast} \\ \rightarrow \text{slow} \end{matrix}$

Dispersion relation for Alfvén wave with perturbed density $\rho = \rho_0 + \tilde{\rho}$:

$$\omega = k_{\parallel} v_A = k_{\parallel} \frac{B_0}{\sqrt{4\pi(\rho_0 + \tilde{\rho})}} \cong k_{\parallel} v_A - \frac{k_{\parallel} v_A}{2} \frac{\tilde{\rho}}{\rho_0}, \quad v_A = \frac{B_0}{\sqrt{4\pi\rho_0}}$$

\downarrow
fast

\downarrow
slow

$$\Rightarrow \omega = \omega^{(0)} + i \left(\frac{\partial}{\partial t} \right)_{\text{slow}}, \quad k_{\parallel} = k_{\parallel}^{(0)} - i \left(\frac{\partial}{\partial x} \right)_{\text{slow}}$$

$$\Rightarrow \frac{\partial}{\partial t} \delta B = - \frac{v_A}{2} \left[\frac{\tilde{\rho}}{\rho_0} \right] \left(\frac{\partial}{\partial x} \delta B \right), \quad x, t \text{ are slow variables.}$$

Envelope eqn \uparrow set by \tilde{v}_{\parallel} dynamics

$$\frac{\partial}{\partial t} \tilde{\rho} = - \rho_0 \frac{\partial}{\partial x} \tilde{v}_{\parallel}$$

$$\rho_0 \frac{\partial}{\partial t} \tilde{v}_{\parallel} = - \frac{\partial}{\partial x} \tilde{\rho} - \frac{\partial}{\partial x} \left(\rho_0 \frac{\tilde{v}_{\parallel}^2}{2} + \frac{\delta B^2}{8\pi} \right)$$

For Alfvén wave, $\rho_0 \frac{\tilde{v}_{\parallel}^2}{2} = \frac{\delta B^2}{8\pi}$
and $\tilde{\rho} = c_s^2 \tilde{v}_{\parallel}$, $c_s^2 = \gamma \frac{\tilde{\rho}}{\rho}$

$$\Rightarrow \frac{\partial^2}{\partial t^2} \tilde{v}_{\parallel} = c_s^2 \frac{\partial^2}{\partial x^2} \tilde{v}_{\parallel} - \frac{\partial^2}{\partial t \partial x} \frac{\delta B^2}{4\pi\rho_0}$$

Transfer to the frame of wave moving, $\tilde{v}_{\parallel} = \tilde{v}_{\parallel}(x - v_A t)$

$$\tilde{v}_{\parallel} = \frac{v_A}{v_A^2 - c_s^2} \frac{\delta B^2}{4\pi\rho_0}, \quad \frac{\tilde{\rho}}{\rho_0} = \frac{\tilde{v}_{\parallel}}{v_A} = \frac{1}{1-\beta} \frac{\delta B^2}{B_0^2}, \quad \beta = \frac{c_s^2}{v_A^2}$$

Plug $\frac{\tilde{\rho}}{\rho_0}$ into the envelope eqn:

$$\frac{\partial}{\partial t} \delta B = - \frac{v_A}{2} \frac{1}{1-\beta} \left[\frac{\delta B^2}{B_0^2} \right] \frac{\partial}{\partial x} \delta B + \left[\begin{matrix} ? & ? & ? \\ \vdots & \vdots & \vdots \end{matrix} \right]$$

steepening v.s. ?? (What stops steepening?)

Steepening is stopped by :

① Dissipation : $\eta \frac{\partial^2}{\partial x^2} \delta B \leftarrow \text{small}$

② Dispersion : $i d_i^2 \Omega_i \frac{\partial^2}{\partial x^2} \delta B \leftarrow \text{ion inertial scale physics.}$

↘ Quasi-parallel Alfvénic collisionless shock

Full envelope eqn :

$$\frac{\partial \delta B}{\partial t} = - \frac{v_A}{z} \frac{1}{1-\beta} \frac{\delta B^2}{B_0^2} \frac{\partial}{\partial x} \delta B + \eta \frac{\partial^2}{\partial x^2} \delta B + i d_i^2 \Omega_i \frac{\partial^2}{\partial x^2} \delta B$$

which is a derivative nonlinear Schrödinger (DNLS) equation.